Will Spin-Relaxation Times in Molecular Magnets Permit Quantum Information Processing?

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Using X-band pulsed electron-spin resonance, we report the intrinsic spin-lattice ($T_1$) and phase-coherence ($T_2$) relaxation times in molecular nanomagnets for the first time. In Cr$_2$M heterometallic wheels, with $M = $ Ni and Mn, phase-coherence relaxation is dominated by the coupling of the electron spin to protons within the molecule. In deuterated samples $T_2$ reaches 3 $\mu$s at low temperatures, which is several orders of magnitude longer than the duration of spin manipulations, satisfying a prerequisite for the deployment of molecular nanomagnets in quantum information applications.

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Certain computational tasks can be efficiently implemented using quantum logic, in which the information-carrying elements are permitted to exist in quantum superpositions [1]. To achieve this in practice, a physical system that is suitable for embodying quantum bits (qubits) must be identified. Some proposed scenarios employ electron spins in the solid state, for example, phosphorous donors in silicon [2], quantum dots [3], heterostructures [4], and endohedral fullerenes [5,6], motivated by the long electron-spin-relaxation times exhibited by these systems. An alternative electron-spin based proposal exploits the large number of quantum states and the nondegenerate configurations in which the moment is “up” or “down” relative to the easy axis. In the absence of a magnetic field, the system spin with a large and negative zero-field splitting can be identified. Some proposed scenarios employ electron manipulations, a prerequisite for the deployment of these systems in quantum information applications.

Using X-band pulsed electron-spin resonance, we report the intrinsic spin-lattice ($T_1$) and phase-coherence ($T_2$) relaxation times in molecular nanomagnets for the first time. In Cr$_2$M heterometallic wheels, with $M = $ Ni and Mn, phase-coherence relaxation is dominated by the coupling of the electron spin to protons within the molecule. In deuterated samples $T_2$ reaches 3 $\mu$s at low temperatures, which is several orders of magnitude longer than the duration of spin manipulations, satisfying a prerequisite for the deployment of molecular nanomagnets in quantum information applications.

Measurements were performed using commercial Bruker Elexsys 580 X-band pulsed ESR spectrometers, employing $^4$He flow cryostats for temperature control. The relaxation times were obtained using standard techniques [26]: $T_2$ from the decay of a 2-pulse Hahn-echo sequence,

$$\pi/2 - \tau - \pi - \tau - \text{echo},$$

with $\tau$ varying; $T_1$ from the recovery of the magnetization (measured with a spin echo) after an inversion pulse,

$$\pi - T - \pi/2 - \pi - \tau - \pi - \tau - \text{echo},$$

with $T$ varying and $\tau$ fixed and short.

Two factors conspire to complicate the measurement of intrinsic lifetimes in anisotropic magnetic clusters. First, in...
a crystal the magnetic cores are typically separated by rather small distances of the order of 1 nm to 10 nm, and are therefore coupled by dipolar interactions. For two free electron spins, the dipolar interaction is of the order of 100/r^3 MHz · nm^3 [26], so this dipolar coupling is the dominant relaxation mechanism limiting T_2 in crystals of clusters. Second, clusters tend to exhibit strongly axial behavior with significant zero-field splittings; this lifts the degeneracy of the \Delta m_s = ± 1 transitions in high spin (S > 1/2) molecules, and leads to a strong dependence of the transition energies on the orientation with respect to the external magnetic field. Thus if we are to study a particular transition in an ensemble of identical molecules, they must be orientationally ordered. The standard approach to solving the first problem, dissolving the crystal (thereby increasing the average separation of the clusters), leads to a second problem, an ensemble of randomly oriented, highly axial molecules.

We chose the two compounds studied here with these factors in mind. Cr_7Ni has a ground state spin of S = 1/2, so it exhibits a single ESR transition and, therefore, no zero-field splitting; the anisotropy of the g factor, which is small, is the only contribution to a dependence of the transition energy on external magnetic field orientation. A dilute dissolved sample of this material is therefore amenable to ESR measurements without causing problems associated with the orientational disorder. Cr_7Mn is a very closely related compound with a ground state spin of S = 1; the similarity in its structure leads us to expect that it should share relaxation mechanisms with Cr_7Ni, but its higher spin allows us to examine the consequences of the zero-field splitting. For example, a modulation of the zero-field splitting, which might occur as a result of coupling with thermally excited mechanical deformations of the molecule, provides further spin-lattice and phase-coherence relaxation mechanisms [27].

Samples were prepared as reported elsewhere [17], and dissolved in toluene. The solution was diluted progressively until the results reported below were no longer dependent on concentration, indicating that the dipolar coupling between clusters had become negligible; this occurred for concentrations below approximately 0.2 mg/ml, corresponding to a mean separation of clusters in excess of about 25 nm.

Figure 1 shows echo-detected ESR as a function of the applied magnetic field measured at 4.5 K for (a) Cr_7Ni and (b) Cr_7Mn. The spectrum for Cr_7Ni shows a single narrow line, as expected for an S = 1/2 species. The line shape is approximately Gaussian, with a width of about 0.01 T, suggesting that the broadening is due to inhomogeneity of the cluster environment. In contrast, the spectrum for Cr_7Mn is broad and contains structure that is characteristic of an S = 1 species with a zero-field splitting exceeding the microwave energy. Also shown is a simulated powder spectrum [28] for an S = 1 species with the Hamiltonian

$$\mathcal{H} = g \mu_B \mathbf{B} \cdot \mathbf{S} + D S_z^2 + E (S_x^2 - S_y^2)$$

(3)

with parameters g = 1.9, D = 21 GHz, and E = 1.9 GHz, which reproduces the main features of the data [green line in Fig. 1(b)]; deviations are probably due to the fact that in the experiment, the pulses deviate from perfect π/2 pulse and 128 ns for a π pulse, ensure that only spins within a window of about 0.3 mT are excited. Using a broad integration window suppresses ^1H ESEEM effects. The echo intensity is proportional to the ESR excitation spectrum. The fine structure in the data close to 0.33 T is an artifact of the cavity.

The strong coupling to protons may provide an efficient phase decoherence path, reducing T_2 [29]. A direct test of

![FIG. 1 (color online). X-band echo-detected ESR as a function of magnetic field for (a) Cr_7Ni and (b) Cr_7Mn measured at 4.5 K (blue), and the simulated powder spectrum for a species with S = 1, g = 2, D = 21 GHz, E = 1.9 GHz (green). In this experiment, the intensity of a Hahn-echo signal with short (τ = 300 ns) delays is measured as a function of the applied magnetic field. Selective pulses, of 64 ns for a π/2 pulse and 128 ns for a π pulse, ensure that only spins within a window of about 0.3 mT are excited. Using a broad integration window suppresses ^1H ESEEM effects. The echo intensity is proportional to the ESR excitation spectrum. The fine structure in the data close to 0.33 T is an artifact of the cavity.](image-url)
FIG. 2 (color online). (a) Decay of the Hahn-echo intensity in Cr$_7$Ni as a function of delay time for long, selective pulses (64 ns $\pi/2$ pulse, 128 ns $\pi$ pulse, blue) and for short broadband pulses (16 ns $\pi/2$ pulse, 32 ns $\pi$ pulse, cyan). (b) Decay of the Hahn-echo intensity for long, selective pulses in perdeuterated Cr$_7$Ni. Dashed green lines indicate fits to the data. The fit in (b) assumes that the ESEEM effect is dominated by a single harmonic at the $^2$D Zeeman frequency [31]. (Note the different horizontal axis scales.)

whether this is indeed an important relaxation mechanism is to measure the same Hahn-echo decay in the perdeuterated analogue compound, as shown in Fig. 2(b). $^2$D has a gyromagnetic ratio about 6 times smaller than $^1$H; the ESEEM frequency is correspondingly about 6 times smaller at 2.556 ± 0.005 MHz, and $T_2$ is about 6 times longer, at 2.210 ± 20 ns. This confirms that the coupling to protons dominates the spin decoherence in the hydrogenated sample. (Note that we were unable to decrease the bandwidth of the pulses sufficiently to suppress the lower frequency ESEEM in the deuterated compound.)

Figure 3 shows the temperature dependences of $T_1$ and $T_2$ for each compound. $T_1$ is comparable between the Cr$_7$Ni and Cr$_7$Mn, and increases rapidly as the temperature falls. This suggests that thermal processes (such as couplings to phonons) are responsible for the longitudinal relaxation. At low temperatures, there is a variation of a factor of about two in $T_1$ measured at different points in the spectrum of Cr$_7$Mn. Exciting different parts of the powder spectrum corresponds to selecting subpopulations of molecules by orientation. These orientational subpopulations relax at different rates, showing that the magnetic anisotropy plays at least some role in the longitudinal relaxation, probably through spectral diffusion effects [30]. There is no such variation of $T_2$.

In each compound, $T_2$ also increases as the temperature is decreased, though less dramatically than $T_1$, and there are signs that it begins to saturate at temperatures below about 2 K. There are two interesting observations: first, there is very little difference between the hydrogenated Cr$_7$Ni and Cr$_7$Mn, despite the strong magnetic anisotropy of the latter; second, the factor of about six between the decoherence times for the perdeuterated and hydrogenated Cr$_7$Ni is retained over the whole temperature range over which a Hahn echo is measurable. Both observations support the hypothesis that dipolar coupling with $^1$H or $^2$D nuclei (which belong to organic ligands of the magnetic cluster and are well-distributed about the core [17]) dominates the phase decoherence in these materials. Phase decoherence arising from fluctuations of the zero-field splitting in Cr$_7$Mn is negligible in comparison over the temperature range studied here.

In conclusion, we have measured intrinsic spin-lattice ($T_1$) relaxation times and, for the first time (to the best of our knowledge), the phase-coherence ($T_2$) relaxation times in molecular magnets. We find that in the heterometalic clusters Cr$_7$Ni and Cr$_7$Mn, $T_1$ is long and somewhat dependent on the magnetic anisotropy of the cluster, but that $T_2$ is dominated by the coupling to the nuclear moments of protons in the vicinity of the cluster. There is no evidence of coupling between the magnetic cluster and the fluorine nuclei, which had previously been identified as a potential decoherence path [9]. Furthermore, we find that the intrinsic phase-coherence time $T_2$ exceeds previous (worst-case) expectations by 3 orders of magnitude, reaching 0.55 $\mu$s.
at 1.8 K, and 3.8 µs for the perdeuterated analogue. With existing apparatus, the time scale for coherent manipulations of the electron spin is of the order of 10 ns; if heteromagnetic clusters of this class were to be used as elements of a quantum information processing device, this would lead to a single-qubit figure of merit of several hundred. The identification of coupling to the 1H or 2D nuclei as the main decoherence path offers a strategy for synthesizing structures with even better coherence properties by reducing as far as possible the number of hydrogens and other magnetic nuclei in the vicinity of the cluster. Our results are very encouraging for the prospects of constructing and manipulating nontrivial quantum states within individual clusters [7,8] and between clusters in composites [11].

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